

Home

Search Collections Journals About Contact us My IOPscience

Comment on 'Irreducible Green function theory for ferromagnets with first- and secondneighbour exchange'

This article has been downloaded from IOPscience. Please scroll down to see the full text article. 1995 J. Phys.: Condens. Matter 7 6967 (http://iopscience.iop.org/0953-8984/7/34/019) View the table of contents for this issue, or go to the journal homepage for more

Download details: IP Address: 171.66.16.151 The article was downloaded on 12/05/2010 at 22:01

Please note that terms and conditions apply.

COMMENT

Comment on 'Irreducible Green function theory for ferromagnets with first- and second-neighbour exchange'

Edward B Brown

Department of Physics, Manhattan College, Riverdale, New York 10471, NY, USA

Received 28 February 1995

Abstract. The 'new' irreducible Green function theory suggested by Mitra and Chakraborty (J. Phys.: Condens. Matter 7 (1995) 379) is shown to be inconsistent and incorrect.

The Heisenberg model Hamiltonian with first- and second-neighbour interactions is given by

$$H = \mu h \sum_{i} S_{i}^{z} - \sum_{i_{1}j_{1}} J_{i_{1}j_{1}}^{(1)} (S_{i_{1}} \cdot S_{j_{1}}) - \sum_{i_{2}j_{2}} J_{i_{2}j_{2}}^{(2)} (S_{i_{2}} \cdot S_{j_{2}})$$
(1)

where $\{S_{i_1}^{\mu}\}$, $\mu = x$, y, z, are the components of a spin operator at site i_1 , $J_{i_1j_1}^{(1)}$ and $J_{i_2j_2}^{(2)}$ are, respectively, the nearest- and next-nearest-neighbour exchange integrals respectively, and the first term represents the coupling of the system to an external z field. The central parameter in the standard irreducible Green function (SIRG) treatment of (1) is the residual momentum space Zubarev Green function (Mitra and Chakraborty 1995),

$$\langle\langle \Phi_{\boldsymbol{k},\boldsymbol{k}'}; S_{\boldsymbol{q}}^{-}\rangle\rangle \equiv \left\langle \left| \left(S_{\boldsymbol{k}'}^{z} S_{\boldsymbol{k}-\boldsymbol{k}'}^{+} - S_{\boldsymbol{k}-\boldsymbol{k}'}^{z} S_{\boldsymbol{k}'}^{+} \right) - C_{mf} - C_{emf}; S_{\boldsymbol{q}}^{-} \right\rangle \right\rangle \tag{2}$$

where C_{mf} and C_{emf} are given by

$$C_{mf} \equiv \langle S_{k'}^z \rangle S_{k-k'}^+ - \langle S_{k-k'}^z \rangle S_{k'}^+ \tag{3}$$

$$C_{emf} \equiv \left(A_{k'} - A_{k-k'}\right)S_k^+. \tag{4}$$

In the SIRG approach, A_k is determined by the irreducibility condition

$$\lambda_{k,k',q} \equiv 0 \tag{5}$$

$$\lambda_{k,k',q} \equiv \langle \left[\Phi_{k-k'}, S_q^- \right] \rangle. \tag{6}$$

Mitra and Chakraborty (1995) have suggested that the irreducibility condition (5), may be replaced by the less restrictive condition that $\lambda_{k,k',q}$ is independent of k'. They have shown that, if $\lambda_{k,k',q}$ is independent of k', it disappears from the relevant Green function equation of motion. They have asserted that the assumed k' independence allows one to choose any A_k (they chose A_k in accordance with the Callen decoupling scheme) since A_k no longer must be chosen to satisfy the irreducibility condition, (5).

0953-8984/95/346967+02\$19.50 © 1995 IOP Publishing Ltd

The purpose of this note is to show that if $\lambda_{k,k',q}$ is independent of k', then it must vanish and A_k must still satisfy the irreducibility condition (5). The proof is straightforward. If $\lambda_{k,k',q}$ is independent of k', it must have the same value for all k'. In particular, we must have

$$\lambda_{k,k',q} = \lambda_{k,k'=0,q} = \lambda_{k,k'=k,q}.$$
(7)

However, from (2), (3) and (4), one readily obtains

$$\lambda_{k,k'=0,q} = -\lambda_{k,k'=k,q} \tag{8}$$

so that (7) gives

$$\lambda_{k,k',q} = 0 \tag{9}$$

if $\lambda_{k,k',q}$ is independent of k'. Thus, A_k must still be chosen to satisfy (9) and, in particular, cannot be chosen in accordance with the Callen decoupling scheme. Therefore, the Mitra-Chakraborty treatment of (1) is inconsistent and incorrect.

References

Mitra S N and Chakraborty K G 1995 J. Phys.: Condens. Matter 7 379